Introduction

In this thesis, we propose a model of a ‘gas of circles’ (GOC), the ensemble of regions in the image domain consisting of an unknown number of circles with approximately fixed radius and short range repulsive interactions. To show its empirical success we apply the model on a forestry application, which is detecting exact tree crowns on remotely sensed images. The method is based on the ‘higher order active contour’ (HOAC) [RJZ03, RJZ06] framework, which incorporates long-range interactions between contour points and thereby include prior geometric information without using a template shape. This makes them ideal when looking for multiple instances of an entity in an image. For such a model to work, the circles must be stable to small perturbations of their boundaries, \textit{i.e.} they must be local minima of the HOAC energy, for otherwise a circle would tend to ‘decay’ into other shapes. This is a nontrivial requirement. We impose it by performing a functional Taylor expansion of the HOAC energy around a circle, and then demanding that the first order term be zero for all perturbations, and that the second order term be positive semi-definite. These conditions allow us to fix one of the model parameters in terms of the others, and constrain the rest. The energy is minimized using gradient descent algorithm, and implemented using the level-set method [OS88]. Experiments using the HOAC energy demonstrate empirically the coherence between these theoretical considerations.

The general ‘gas of circles’ model has many potential applications in varied domains, but it suffers from a drawback: the local minima corresponding to circles can trap the gradient descent algorithm, thus producing phantom circles even with no supporting data. We solve the problem of phantom circles by calculating, via a Taylor expansion of the energy, parameter values that make the circles into inflection points rather than minima. In addition, we find that this constraint reduces the prior parameters one overall weight, while improving the empirical success of the model.

Although, the HOAC ‘gas of circles’ model is an effective tool for modeling circular shapes, there are some difficulties. It is complicated to express the space of regions in the contour representation, and consequently difficult to work with a probabilistic formulation. In addition, from the algorithmic point of view, the current model does not allow enough topological freedom, and the implementation of the HOAC model is difficult and computationally expensive. But, it is possible to create an alternative formulation of HOAC models, based on the ‘phase field’ framework much used in physics to model regions and interfaces. The standard phase field model is, to a good approximation, equivalent to a classical active contour model. We compute, as a function of the HOAC energy parameters, the phase field energy parameters that produce an equivalent model. This means that we can adjust the phase field parameters to ensure stable circles of a given radius also. We extend the phase field ‘gas of circles’ model to the case of an inflection point rather than a minimum in the circle energy at the desired radius.

We also introduce two data models. The first describes the use of only one, the most significant, infrared band of the three available bands in the colour infrared (CIR) images. The model is based on the image gradient and on Gaussian distributions, with the values at different pixels independent, and with different means and variances for tree crowns and the background. While successful, this model, even with the strong region prior, is not capable of extracting accurately the borders of all the trees. Some trees are simply too similar to the background. To solve this problem our second data model makes use of all three bands in the CIR images. We study the improvement or otherwise of the extraction results produced by modelling the three bands as independent or as correlated. As we will see, even at the level of maximum likelihood, the inclusion of ‘colour’ information, and in particular, interband correlations, can improve the results, and in conjunction with the region prior, the full model is considerably better than that based on one band alone.

We use the above models to extract tree crowns on aerial images, but the models are not restricted to forest management. They can also be applied to the detection of other circular objects. A few other examples: in nanotechnology to detect various particles and microarrays in electron-microscopic images; in biology to segment circular cells and molecules, or to detect pollen grains; in medical image processing; in remote sensing, in the processing of aerial and satellite images, for meteorological, military, and agricultural management.
**Thesis 1: The ‘gas of circles’ model [HJKZ06a]**

HOACs include multiple integrals over the contour. These integrals correspond to long-range interactions between tuples of contour points, and allow the incorporation of sophisticated prior geometric knowledge. Combined with length and area terms, one of the basic forms of Euclidean invariant quadratic HOAC models [RJZ03, RJZ06] is

\[
E_g(\gamma) = \lambda_C L(\gamma) + \alpha_C A(\gamma) - \frac{\beta_C}{2} \int \int dp \ dp' \ \dot{\gamma}(p) \cdot \dot{\gamma}(p') \ \Psi(R(p, p')) ,
\]

where \( \gamma \) is the contour, parameterized by \( p \); \( L \) is the length of the contour; \( A \) is its interior area; \( R(p, p') = |r(p, p')| \), where \( r(p, p') = \gamma(p) - \gamma(p') \); and \( \Psi \) is an interaction function that determines the geometric content of the model. Via the stability analysis we adjust the parameters of this model so that the low energy configurations are approximate circles of approximately fixed radius.

**My contribution**

*The Author presented the stability analysis of the higher-order active contour model for a ‘gas of circles’, defined the required conditions and consequent parameter constraints.*

**Stability analysis**

The model is of the form \( E(\gamma, I) = E_i(I, \gamma) + E_g(\gamma) \). The likelihood energy \( E_i \) is described in the 4\(^{th} \) thesis. We analyze the prior energy \( E_g \), and show how it can be used to model a ‘gas of circles’. We adjust the parameters of the model so that configurations consisting of collections of circles of approximately a certain radius \( r_0 \) are stable, and have low energy. Stability means that, if the shape of a circle of radius \( r_0 \) is changed slightly, it will relax back into the circle. We choose parameters so that a circle of radius \( r_0 \) is a minimum of \( E_g \). We expand the energy in a Taylor series to second order in perturbations around a circle of radius \( r_0 \), then adjust the parameters so that the first derivative is zero, which tells that the circle is an energy extremum, and so that the second derivative is positive definite, so the extremum is a minimum. The parameters can be further adjusted so that the energy of the minimizing circle is not too high.

We calculate \( E_g(\gamma) = E_g(\gamma_0 + \delta \gamma) \) to second order in \( \delta \gamma \), where \( \gamma_0 \) is a circle of radius \( r_0 \). Since we are expanding around a circle, it is easiest to use polar coordinates \((r, \theta)\), and to choose \( \theta(p) = p \) as the parameterization. Tangential changes \( \delta \theta \) can be undone by a diffeomorphism and hence do not affect the energy. The radial perturbation \( \delta r \) can be expanded in a Fourier series on the circle:

\[
\delta r(p) = \sum_k a_k e^{ikr_0p} \quad k = m/r_0, \ m \in \mathbb{Z} ,
\]

where \( a_k \ll r_0 \). The \( L(\gamma) \) and the \( A(\gamma) \) can be expressed to second order as:

\[
L(\gamma) = 2\pi r_0 \left\{ 1 + \frac{a_0}{r_0} + \frac{1}{2} \sum_k k^2 |a_k|^2 \right\} \quad (2a)
\]

\[
A(\gamma) = \pi r_0^2 + 2\pi r_0 a_0 + \pi \sum_k |a_k|^2 . \quad (2b)
\]

The expansion of the quadratic term in equation (1) is more complicated, since one have to expand \( \dot{\gamma}, R, \) and \( \Psi \), but invariance with respect to translations around the circle means that the second-order term is
diagonal in the Fourier basis. The result, after combination with equations (2), is that the prior energy is given to second-order by:

\[ E_g(\gamma_0 + \delta \gamma) = E_0 + a_0 E_1 + \frac{1}{2} \sum_k |a_k|^2 E_2(k) , \]

where

\[
\begin{align*}
E_0 &= 2\pi \lambda_C r_0 + \pi \alpha_C r_0^2 - \pi \beta_C \int_0^{2\pi} dp \text{ } F_{00} \\
E_1 &= 2\pi \lambda_C + 2\pi \alpha_C r_0 - 2\pi \beta_C \int_0^{2\pi} dp \text{ } F_{10} \\
E_2 &= 2\pi \lambda_C r_0 k^2 + 2\pi \alpha_C - 2\pi \beta_C \left[ \left( 2 \int_0^{2\pi} dp \text{ } F_{20} + \int_0^{2\pi} dp \text{ } F_{21} e^{ikr_0p} \right) \\
&\quad + k \left( 2ir_0 \int_0^{2\pi} dp \text{ } F_{23} e^{ikr_0p} \right) \right] + k^2 \left( r_0^2 \int_0^{2\pi} dp \text{ } F_{24} e^{ikr_0p} \right) .
\end{align*}
\]

(3)

The \( F \) are functions of \( p \) and \( r_0 \), and depend on \( \Psi \).

**Parameter constraints**

In order to achieve an extremum, the linear term must be zero. This implies that:

\[ \beta_C(\lambda_C, \alpha_C, r_0) = \frac{\lambda_C + \alpha_C r_0}{\int_0^{2\pi} dp \text{ } F_{10}} , \]

(4)

which fixes \( \beta_C \) given \( \lambda_C \), \( \alpha_C \), and \( r_0 \). One can set \( \lambda_C = 1.0 \) without loss of generality. Given \( r_0 \), we therefore have only two free parameters to adjust to achieve stability, i.e. to make \( E_2 \) positive for all \( k \). In figure 1 left, can be seen the plot of \( \beta_C \) versus \( r_0 \) and \( \alpha_C \). A given \( r_0 \) defines a slice of the surface, the potentially stable pairs of \( (\alpha_C, \beta_C) \), but only those pairs for which \( E_2 \geq 0 \) for all \( k \) are actually stable. Figure 1 right shows the plot of \( E_0 \) for one such pair. Note that it has a minimum at \( r_0 \), so that it is stable against radial perturbations \( (k = 0) \). It is also stable with respect to perturbations with \( k \neq 0 \).

Figure 2 shows experiments using just the prior energy, starting from various initial conditions, illustrating the formation of circles of the desired radius.
Figure 2: Formation of stable circles, second and fourth, with $r_0 = 10$ and 20 respectively, from two different initial conditions, first and third.

**Thesis 2: The inflection point ‘gas of circles’ model [HJKZ06b]**

Figure 1 right shows a plot of the energy of a circle versus radius for parameter values selected according to the ‘gas of circles’ criteria, it contains a local energy minimum at $r_0$. In the absence of supporting data, the global minimum will be the empty region, the correct behaviour. A gradient descent algorithm, however, cannot escape from these local minima, meaning that circles of radius $r_0$, once formed during gradient descent, cannot disappear, even if the data does not support their existence.

**My contribution**

The Author presented new parameter settings of the HOAC energy that creates an inflection point in the energy function rather than a minimum, so that circles vanish without supporting image data, but a small amount of image information can create an energy minimum at the given radius. This criteria fixes all parameters of the prior except one.

**Inflection point**

The idea we pursue is to adjust the parameters so that the minimum of the curve is replaced by a broad, approximately flat area, as shown in the three plots in figure 4. Such an energy means that in the absence of image data, a circle will shrink and disappear, whereas small amounts of image data will be sufficient to create a minimum in the flat area, thus producing a stable circle. The natural method to achieve such a broad flat region is to create an energy function that has a single inflection point. It is a nontrivial exercise to find parameter values that result in inflection points.

We still require that a circle of radius $r_0$ be stable to sinusoidal perturbations with $k > 0$, but now we also require that such a circle be an inflection point with respect to perturbations with $k = 0$, that is, changes of radius. We will see that these demands are sufficient to fix the prior energy $E_g$ up to an overall multiplicative constant and a small range of values for $d$. More precisely, we still require that $E_1(r_0) = 0$ and $E_2(k, r_0) > 0$ for $k > 0$, but we now require that $E_2(0, r_0) = 0$ too. The first condition gives equation (4). The second condition, which follows from the last term of equation (3), also relates $\alpha_C$ and $\beta_C$:

$$\alpha_C(r_0) = \beta_C(r_0) \tilde{G}(r_0),$$

(5)

where $\tilde{G}(r_0) = 2G_{20}(r_0) + G_{21}(0, r_0)$. We can solve equations (4) and (5) for $\alpha_C$ and $\beta_C$, giving

$$\alpha_C(r_0) = \frac{\lambda_C \tilde{G}(r_0)}{G_{10}(r_0) - r_0 \tilde{G}(r_0)}$$

and

$$\beta_C(r_0) = \frac{\lambda_C}{G_{10}(r_0) - r_0 \tilde{G}(r_0)}.$$
These equations fix $\alpha_C$ and $\beta_C$ as functions of $r_0$ and $d$. Since $r_0$ is fixed by the application, the only remaining parametric degrees of freedom are the value of $d$, and the overall strength of the prior term, represented by $\lambda_C$. Recall, however, that we also require $\alpha_C$ and $\beta_C$ to be positive. The question is then how to find values of $d$ for a given $r_0$ so that $\alpha_C(r_0) > 0$ and $\beta_C(r_0) > 0$.

**Determination of $d$**

To illustrate the behaviour we want to understand, figure 3 shows plots of $\alpha_C$ and $\beta_C$ against $d$ for fixed $r_0$, in this case $r_0 = 5$. There are two critical points, $d_{\min}$ and $d_{\max}$. Only for the range $d_{\min} < d < d_{\max}$ are both $\alpha_C$ and $\beta_C$ positive. Our goal is therefore to find $d_{\min}$ and $d_{\max}$ as functions of $r_0$.

From equations (6), it can be seen that $d_{\max}$ arises from a zero in the denominator, while $d_{\min}$ arises from a zero in the numerator. It is therefore sufficient to find these zeros in order to find $d_{\min}$ and $d_{\max}$. To proceed, we note a scaling property of $G_{00}$. The function $G_{00}$ is given by the following integral:

$$G_{00}(r) = \int_{-\pi}^{\pi} dp \cos(p) r^2 \Psi\left(2\sin\frac{p}{2}\right).$$

Since $\Psi(z)$ is a function of $z/d$ only, by pulling $d^2$ out of the integral one can write $G_{00}$ as $G_{00}(r) = d^2 \hat{G}_{00}(r/d)$. Now recall that $G_{10} = \frac{1}{2} \partial_c G_{00}$ and $\hat{G} = \partial_c \hat{G}_{10}$. Then find that:

$$\hat{G}(r_0) = \hat{G}(r_0/d) \quad \text{and} \quad G_{10}(r_0) - r_0 \hat{G}(r_0) = d\left(\hat{G}_{10}(r_0/d) - \frac{r_0}{d} \hat{G}(r_0/d)\right),$$

where $\hat{G}_{10}(z) = \frac{1}{2} \partial_z \hat{G}_{00}(z)$ and $\hat{G}(z) = \partial_z \hat{G}_{10}(z)$. Thus both numerator and denominator of equations (6) can be written, up to multiplication by positive coefficients, as functions of $r_0/d$. Now, $f(r, d) = \hat{f}(r/d)$ and $f(r, d_0) = 0$ imply $f(ar, ad_0) = 0$ for all $a \in \mathbb{R}$; thus if we determine $d_{\min}$ and $d_{\max}$ for one value of $r_0$, their values are known for any $r_0$.

To determine $d_{\min}$ and $d_{\max}$ while avoiding iterative numerical procedures to find these points, we use a polynomial approximation to $G_{00}$:

$$G_{00}(r) = \sum_{n=0}^{\infty} b_n r^n.$$

It is easy to show that:

$$b_m = \begin{cases} 
0 & m < 2, \\
\frac{1}{(m-2)!} \int_{-\pi}^{\pi} dp \cos(p) Y^{(m-2)}(0) & m \geq 2,
\end{cases}$$

where $Y(r) = \Psi(2r \sin(p/2))$. The derivatives of $Y$ evaluated at zero are:

$$\frac{Y^{(m)}(0)}{(2|\sin(p/2)|)^m} = \Psi^{(m)}(0) = \begin{cases} 
1 & m = 0, \\
0 & m = 1 \text{ or } m \text{ even}, \\
(-1)^{\frac{m-1}{2}} \left(\frac{\pi}{2}\right)^{m-1} & m \geq 3 \text{ and } m \text{ odd}.
\end{cases}$$
The contour parameters are given by [RJZ05]:

\[ \text{interface of finite width around } \partial R \]

The phase field model is approximately equivalent to a classical active contour [RJZ05] in the sense that:

\[ \arg \min_{\phi : \zeta(x) = R} E_0(\phi), \text{ i.e. the minimizing phase field for a given fixed region } R, \] would take the value 1 inside \( R \) and −1 outside. The effect of \( D \neq 0 \) is to smooth \( \phi_R \) so that it has an interface of finite width around \( \partial R \). The phase field model is approximately equivalent to a classical active contour [RJZ05] in the sense that:

\[ E_0(\phi_R) \approx \lambda_C L(\partial R) + \alpha_C A(R), \]

The contour parameters are given by [RJZ05]:

\[ \alpha_C = 4\alpha/3, \quad \lambda_C^2 = 16D\lambda K/15, \quad K = 1 + 5(\alpha/\lambda)^2, \quad w^2 = 15D/\lambda K. \]

Figure 4: Plot of \( E_0 \) against \( r \) for \( r_0 = 5.0 \) with \( \alpha \) and \( \beta \) determined by equations (6). First, \( d = 6.4 \); second, \( d = 6.8 \); third, \( d = 7.2 \). For this value of \( r_0 \), \( d_{\min} = 6.3880 \), \( d_{\max} = 7.2495 \).

Substituting into equation (8) gives \( b_m \):

\[
 b_m = \begin{cases} 
 0 & \text{if } m < 5 \text{ or } m \text{ even}, \\
 (-1)^{\frac{m-1}{2}} \frac{4(2\pi)^{m-3}}{m!(m-3)!!} \frac{1}{d^{m-2}} & \text{if } m \geq 5 \text{ and } m \text{ odd}.
\end{cases}
\]

One can then derive expressions for \( \tilde{G} \) and \( G_{10} - r\tilde{G} \):

\[
 \tilde{G}(r) = 2 \sum_{m \geq 3} \frac{(-1)^{\frac{m-1}{2}} (2\pi)^{m-1}(m+1)}{m!!(m-2)!!} \left( \frac{r}{d} \right)^m
\]

\[
 G_{10}(r) - r\tilde{G}(r) = 2d \sum_{m \geq 4} \frac{(-1)^{\frac{m-2}{2}} (2\pi)^{m-2}}{[(m-3)!!]^2} \left( \frac{r}{d} \right)^m.
\]

We computed the roots of these polynomials including terms up to \( m = 49 \). The smallest positive roots furnish the values of \( d_{\min} \) and \( d_{\max} \). The result is that \( d_{\min} \approx 1.2776r_0 \) and \( d_{\max} \approx 1.4499r_0 \). The graphs in figure 4 show plots of \( E_0 \) against \( r \) for \( r_0 = 5 \), with \( d \) values chosen from the domain \( d_{\min} < d < d_{\max} \).

**Thesis 3: The phase field ‘gas of circles’ model [HJ07a, HJ07b]**

A phase field \( \phi \) is a real-valued function on the image domain \( \Omega \). Given a threshold \( z \), a phase field determines a region by the map \( \zeta_\phi(x) = \{ x : \phi(x) > z \} \). Thus phase fields are a level set representation. The difference with the usual distance function level set representation is that the functions are not constrained: the set of possible \( \phi, \Phi \), is a linear space. The simplest phase field energy is [RJZ05]:

\[
 E_0(\phi) = \int_{\Omega} dx \left\{ \frac{D}{2} \partial_x \partial_y \phi + \lambda \left( \frac{1}{4} \phi^4 - \frac{1}{2} \phi^2 \right) + \alpha \left( \phi - \frac{1}{3} \phi^3 \right) \right\}. 
\]

With \( D = 0, \phi_R \triangleq \arg \min_{\phi : \zeta(\phi) = R} E_0(\phi), \) i.e. the minimizing phase field for a given fixed region \( R \), would take the value 1 inside \( R \) and −1 outside. The effect of \( D \neq 0 \) is to smooth \( \phi_R \) so that it has an interface of finite width around \( \partial R \).
To approximate the quadratic HOAC energy, following nonlocal term was added to $E_0$ by Rochery et al. [RJZ05]:

$$E_{NL}(\phi) = -\frac{\beta}{2} \int_{\Omega^2} dx \, dx' \, \partial \phi(x) \cdot \partial \phi(x') \, G(x - x'),$$

(11)

where $G(x - x') = \Psi(|x - x'|/d)$. With $\beta_C = 4\beta$, $E_g = E_0 + E_{NL}$ is equivalent (as usual, to a good approximation) to the HOAC model, and can be used in its place, thus allowing the incorporation of non-trivial prior knowledge about region geometry while still profiting from all the advantages of the phase field framework.

**My contribution**

The Author showed how to transform the previously determined ‘gas of circles’ parameters into the phase field model so that it creates stable circles with the desired radius and determined the constraints on the parameter conversion. The Author presented the steps of the transformation and the constraints converting the inflection point ‘gas of circles’ parameters to the phase field framework.

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**The ‘gas of circles’ higher-order phase field model**

Using the relations between the phase field and contour parameters, we create a phase field ‘gas of circles’ model equivalent to the HOAC model described in Thesis 1 (units are chosen so that $d = 1$, $\tilde{\cdot}$ denotes the parameters normalized by $\lambda_C$):

1. Choose $w$. It cannot be too small, or a subpixel discretization will be needed, and it cannot be too large or the phase field model will not be a good approximation to the HOAC model.

2. Choose $\tilde{\alpha}_C \leq \sqrt{5} / (2w)$.

3. Determine the $\tilde{\beta}_C$ parameter corresponding to $r_0$ and $\tilde{\alpha}_C$ using the method defined in Thesis 1.

4. Set $\tilde{\lambda} = \frac{15}{8w}[1 + \sqrt{1 - 4\tilde{\alpha}_C^2w^2/5}]$, $\tilde{\alpha} = 3\tilde{\alpha}_C/4$, $\tilde{\beta} = \tilde{\beta}_C/4$, and $\tilde{D} = w/4$.

5. Choose $\lambda_C$ and multiply $\tilde{D}$, $\tilde{\lambda}$, $\tilde{\alpha}$, and $\tilde{\beta}$ by it to get $D$, $\lambda$, $\alpha$, and $\beta$.

**The phase field inflection point ‘gas of circles’ model**

We combine the phase field ‘gas of circles’ model with the constraint that the circle energy have an inflection point rather than a minimum at the desired radius. This is a non-trivial requirement. The relations between the phase field and contour parameters were derived using an approximate ansatz for $\phi_R$ by Rochery et al. [RJZ05]. For the ‘gas of circles’ model, the approximations are not expected to be important, since small errors in the parameters will produce small changes in behaviour. However, it is important to see whether these approximate parameter relations preserve the inflection point behaviour.

In the previous subsection, we chose $\tilde{\alpha}_C$ according to the constraint $\tilde{\alpha}_C \leq \sqrt{5} / (2w)$. In Thesis 2, we described how a given value of $r_0$ and $d$ fixes $\tilde{\alpha}_C$ and $\tilde{\beta}_C$. To satisfy both these constraints one need to choose $w \leq \sqrt{5} / 2\tilde{\alpha}_C$. To determine the parameters of the new phase field inflection point ‘gas of circles’ model, we therefore take the following steps:

1. Choose a $d$ value satisfying the inflection point criteria given in Thesis 2. This fixes $\tilde{\alpha}_C$ and $\tilde{\beta}_C$.

2. Choose $w$ using the above criterion.
3. Determine $\tilde{\lambda}$, $\tilde{\alpha}$, $\tilde{\beta}$, and $\tilde{D}$ as before.

4. Multiply these parameters by $\lambda_C$.

**Thesis 4: Data terms for HOAC and phase field models [HJKZ06a, HJ07a, Hor07]**

**My contribution**

The Author introduced two different data models. First, the HOAC model was completed with a term based on the image gradient and Gaussian-based intensity terms coming from a single-spectral image. Second, the previous intensity term was extended to multispectral images using a multi-dimensional Gaussian model.

The simplest likelihood energy, uses single spectral images. We model the image in $R$, and in the background $\bar{R}$, using Gaussian distributions. We add a term that predicts high gradients along the boundary $\partial R$:

$$E_{C,1}(I, R) = \lambda_i \int dp \ n(p) \cdot \partial I(\gamma(p)) + \alpha_i \left[ \int_R d^2 x \ \frac{(I(x) - \mu_{in})^2}{2\sigma_{in}^2} + \int_{\bar{R}} d^2 x \ \frac{(I(x) - \mu_{out})^2}{2\sigma_{out}^2} \right],$$

where $n$ is the (unnormalized) outward facing normal. Note that to facilitate comparison of parameters in the prior energy, we set $\lambda = 1$ in $E_g$ and introduce a weight $\alpha_i$ in $E_i$. The parameters $\mu_{in}$, $\sigma_{in}$, $\mu_{out}$, and $\sigma_{out}$ are learned from examples using maximum likelihood, and then fixed.

The equivalent phase field data term can be defined as:

$$E_i(I, R) = \int_{\Omega} dx \ \left\{ \lambda_i \partial I \cdot \partial \phi + \alpha_i \left[ \frac{(I - \mu_{in})^2}{2\sigma_{in}^2} \phi_+ + \frac{(I - \mu_{out})^2}{2\sigma_{out}^2} \phi_- \right] \right\},$$

where $\phi_\pm = (1 \pm \phi)/2$.

In some of the cases not only single-spectral images but color or multispectral images are also available. To use the opportunity of the other channels we developed a multispectral data term for the HOAC GOC model based on Gaussian distributions with full covariance matrices. The parameters of $E_i$ are learnt from samples of each class using maximum likelihood, and then fixed. We denote the mean vectors inside and outside as $M_{in}$ and $M_{out}$ and the covariance matrices $\Sigma_{in}$ and $\Sigma_{out}$. We define the energy as:

$$E_{C,1}(I, R) = -\int_R dp \ \ln \left[ \det^{-1/2}(\Sigma_{in}/2\pi)e^{-\frac{1}{2}(I(p)-M_{in})^T\Sigma_{in}^{-1}(I(p)-M_{in})} \right]$$

$$-\int_{\bar{R}} dp \ \ln \left[ \det^{-1/2}(\Sigma_{out}/2\pi)e^{-\frac{1}{2}(I(p)-M_{out})^T\Sigma_{out}^{-1}(I(p)-M_{out})} \right].$$

The energy is minimized using gradient descent.

**Experimental results**

In this section we present experimental results were obtained on aerial images provided by the Hungarian Central Agricultural Office, Forestry Administration (CAO, FA) and French Forest Inventory (IFN).
Figure 5: From left to right: image of poplars ©IFN; the best result with a classical active contour; result with the ‘gas of circles’ model.

exact parameter values and the details of the segmentation can be found in the experimental results part of the thesis.

In figure 5 the data is on the left, the best result obtained with a classical active contour model is in the middle, and the result with the HOAC ‘gas of circles’ model is on the right. The trees are close together, using the classical active contour, the result is that the tree crown boundaries touch in the majority of cases, despite their separation in the image. The HOAC model produces more clearly delineated tree crowns, but there are still some joined trees.

Figure 6: From left to right: an image of regularly planted poplars with different fields in the right part ©IFN; result with the phase field ‘gas of circles’ model; result with the inflection point GOC model.

Figure 6 left shows a difficult image to analyse. It has two fields with different intensities on the right. The result with the phase field ‘gas of circles’ model is shown in figure 6 middle. This result clearly demonstrate the disadvantage of the non-inflection point model: phantom objects are created in the homogenous areas. Figure 6 right shows the result with the inflection point ‘gas of circles’ phase field model. The result is very good, with only one false positive. Both results were obtained in less than 1 minute.

Finally figure 7 left shows a different type of CIR image, of isolated trees in fields. The result with the multispectral GOC model, shown in figure 7 middle, is correct, ignoring the field, for example. The result with the model in Thesis 2, in the right image, is not as good, with one large false positive, and smaller errors on each of the detected trees.
Summary of the Author’s contributions

We summarize the results of the Author arranging them into four theses. Table 1 illustrates how the thesis points are covered by the publications of the Author.

1. The Author presented the stability analysis of the higher-order active contour model for a ‘gas of circles’, defined the required conditions and consequent parameter constraints.

2. The Author presented new parameter settings of the HOAC energy that creates an inflection point in the energy function rather than a minimum, so that circles vanish without supporting image data, but a small amount of image information can create an energy minimum at the given radius. This criteria fixes all parameters of the prior except one.

3. The Author showed how to transform the previously determined ‘gas of circles’ parameters into the phase field model so that it creates stable circles with the desired radius and determined the constraints on the parameter conversion. The Author presented the steps of the transformation and the constraints converting the inflection point ‘gas of circles’ parameters to the phase field framework.

4. The Author introduced two different data models. First, the HOAC model was completed with a term based on the image gradient and Gaussian-based intensity terms coming from a single-spectral image. Second, the previous intensity term was extended to multispectral images using a multi-dimensional Gaussian model.
Publications of the Author


Bibliography


Coauthor’s declaration

June 27, 2007

I hereby certify that I am familiar with the following publications of the applicant Mr Péter Horváth:


Regarding to our joint results referred to in these articles, the following ones were obtained as the result of joint contribution by the applicant and myself:

1. The introduction and state-of-the-art. Conclusion.

2. The Introduction and the Conclusion. Fixing the parameters giving an inflection point. Determination of the interaction function parameter providing non-negative parameters.
3. Using the recent inflection point ‘gas of circles’ model parameters with the phase field framework.

4. The application of higher-order phase field model instead of the higher-order active contour model, for the description of multiple circles with similar radii.

The applicant’s contribution was prominent in obtaining the following results:

1. Stability analysis and conditions, and consequent parameter constraints. Addition of the gradient term to data energy. The synthetic justification of the model and the experiments and analysis of results on aerial images.

2. The idea to use an inflection point instead of a minimum in the circle energy, thereby solving the problem of phantom circles. The experiments and analysis of results.

3. Converting the inflection point ‘gas of circles’ parameters, to the phase field framework, determining the steps of the transformation and the constraints. The synthetic and real experiments and analysis of results.

4. Determining the steps and the constraints of the parameter conversion from the higher-order active contours to the phase field framework. Synthetic and real experiments and analysis of results.

My own contribution was prominent in obtaining the following results:

1. Using the higher-order active contour model to describe circular regions. The introduction of the data term using the pixel intensities.

2. Using Taylor polynomials to approximate the interaction function’s parameter.

3. –

4. Embedding recent higher-order active contour data model into the phase field framework.


Dr. Ian H. Jermyn
Coauthor’s declaration

July 31, 2007

I hereby certify that I am familiar with the following publications of the applicant Mr Péter Horváth:


Regarding to our joint results referred to in these articles, the following ones were obtained as the result of joint contribution by the applicant and myself:

1. publication: The introduction and overview on the state-of-the-art. Conclusion.

2. publication The introductory and the conclusion part. Fixing the parameters so that using the derivative constraints giving inflection point. Determination of the interaction function parameter providing non-negative parameters.

The applicant’s contribution was prominent in obtaining the following results:

1. publication: The expression of the quadratic term, stability conditions. The synthetic justification of the model and the experimental results on aerial images.
2. publication: The idea using an inflection point instead of energy minimum at the circle energy, solving the problem of phantom circles. The experimental results part.

My own contribution was prominent in obtaining the following results:

1. publication: Tests on real data.

2. publication: Validation of the experimental results.


Zoltán Kató
Coauthor’s declaration

June 27, 2007

I hereby certify that I am familiar with the following publications of the applicant Mr Péter Horváth:


Regarding to our joint results referred to in these articles, the following ones were obtained as the result of joint contribution by the applicant and myself:

1. The introduction and state-of the art. Conclusion.

2. The Introduction and the Conclusion. Fixing the parameters giving an inflection point. Determination of the interaction function parameter providing non-negative parameters.

The applicant’s contribution was prominent in obtaining the following results:

1. Stability analysis and conditions, and consequent parameter constraints. Addition of the gradient term to data energy. The synthetic justification of the model and the experiments and analysis of results on aerial images.

2. The idea to use an inflection point instead of a minimum in the circle energy, thereby solving the problem of phantom circles. The experiments and analysis of results.
My own contribution was prominent in obtaining the following results:

1. Test on real data
2. Validation with the French Food Institute


Dr. Josiane Zerubia